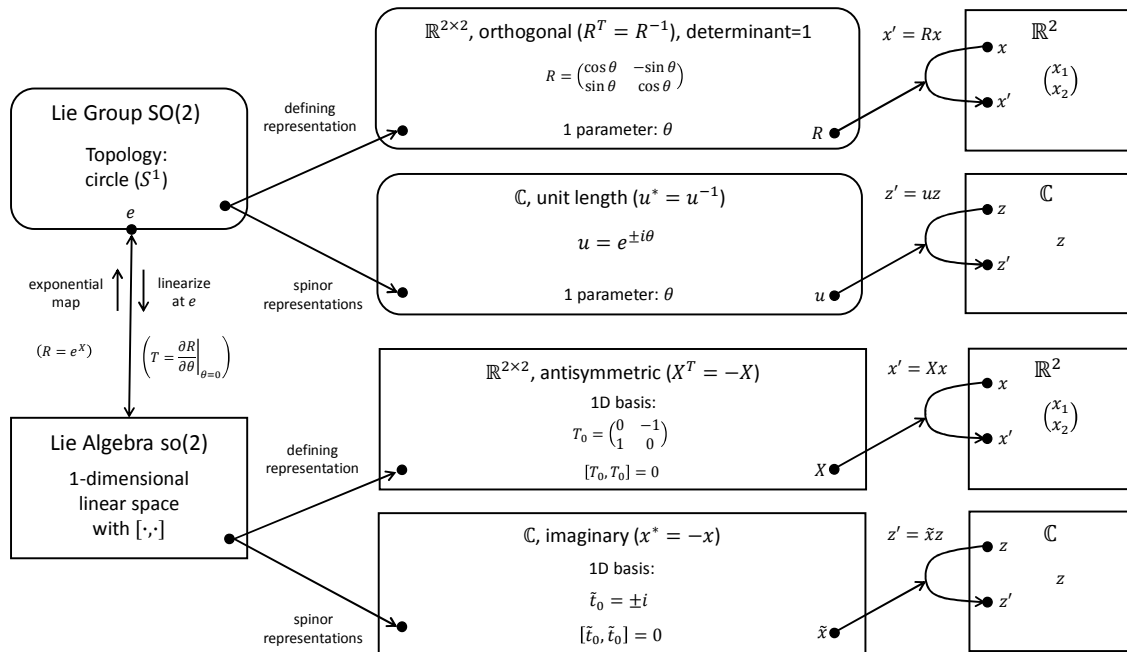


6.13 SO(2): Spinor Representations; Spin(2) = SO(2)



In the previous example, we showed that $SO(2)$ and $U(1)$ are isomorphic. Hence, all representations of $U(1)$ are also representations of $SO(2)$. This includes the 1-dimensional representations of the form $u = e^{ik\theta}$, where $k = 0, \pm 1, \pm 2, \dots$. But these representations act on a *complex* variable: they are *spinor representations* of $SO(2)$! The lowest two spinor representations (with $k = \pm 1$), usually designated S^+ and S^- , are shown in the lower branch of the diagram [GTNut, Ch. VII.1, p. 410]. For comparison, the defining (2-dimensional) representation of $SO(2)$ is shown again in the upper branch.

These 1-dimensional spinor representations are the *irreducible representations* of $SO(2)$. Remember that the irreducible representations of a commuting (= Abelian) group, such as $SO(2)$, must be one dimensional [QTGR, Ch. 2.1].

We know from our discussion of $U(1)$ that the S^+ and S^- representations are inequivalent, that is, there is no similarity transformation taking one to the other. Moreover, they have opposite handedness or *chirality* because they are related by the substitution $\theta \rightarrow -\theta$.

We first encountered spinor representations when discussing $SO(3)$. There, the $so(3)$ algebra had even-dimensional spinor representations, with the lowest one being two dimensional. Furthermore, when discussing $SO(4)$ and later $SO^+(1,3)$, we found that the algebra had spinor representations, with the lowest ones being the 2-dimensional representations $(\frac{1}{2}, 0)$ and $(0, \frac{1}{2})$. These representations were irreducible, inequivalent, and had opposite chirality, just like the S^+ and S^- representations of $SO(2)$.

But there is something funny going on: Exponentiating the spinor representations of $so(3)$ led us to representations of $Spin(3)$, the double cover of $SO(3)$. Similarly, exponentiating the spinor representations of $so(4)$ led us to representations of $Spin(4)$, the double cover of $SO(4)$. So, we would expect that exponentiating the spinor representations of $so(2)$ takes us to representations of $Spin(2)$,

the double cover of $SO(2)$. But the $e^{\pm i\theta}$ transformations, shown in the diagram, do *not* double cover $SO(2)$: a 360° rotation doesn't take z to $-z$.

It turns out that $Spin(2)$ and $SO(2)$ are isomorphic [nLab, ncatlab.org]! Topologically, both of them are circles. So, taking the exponential map of $so(2)$ does take us to $Spin(2)$, but it's *not* a new group, it's just $SO(2)$ again. (In contrast, the *universal* cover of $SO(2)$ is a new group, namely \mathbb{R} .)

For $SO(4)$ and $SO^+(1,3)$ we were able to construct a nonchiral spinor representation by taking the direct sum of its left- and right-chiral (half) spinor representations: $(\frac{1}{2}, 0) \oplus (0, \frac{1}{2})$. Analogously, we can construct a nonchiral spinor representation of $SO(2)$ by taking the direct sum of its two chiral (half) spinor representations, $S^+ \oplus S^-$:

$$\begin{pmatrix} z'_1 \\ z'_2 \end{pmatrix} = \begin{pmatrix} e^{i\theta} & 0 \\ 0 & e^{-i\theta} \end{pmatrix} \cdot \begin{pmatrix} z_1 \\ z_2 \end{pmatrix}.$$

Flipping the chirality of this representation by substituting $\theta \rightarrow -\theta$ swaps the matrix components on the diagonal. Now, this is a similarity transformation, $U \rightarrow SUS^{-1}$,

$$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \cdot \begin{pmatrix} e^{i\theta} & 0 \\ 0 & e^{-i\theta} \end{pmatrix} \cdot \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} e^{-i\theta} & 0 \\ 0 & e^{i\theta} \end{pmatrix}$$

and therefore this 2-dimensional spinor representation is nonchiral. Moreover, it is equivalent to the defining representation of $SO(2)$:

$$\frac{1}{\sqrt{2}} \begin{pmatrix} i & -i \\ 1 & 1 \end{pmatrix} \cdot \begin{pmatrix} e^{i\theta} & 0 \\ 0 & e^{-i\theta} \end{pmatrix} \cdot \frac{1}{\sqrt{2}} \begin{pmatrix} -i & 1 \\ i & 1 \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}.$$

Note how different this is from spinors in four dimensions: the 4-dimensional Dirac spinor representation and the defining 4-vector representation were very different animals!

For another perspective on the spinor representation of $SO(2)$, let's construct it by starting from the Clifford algebra $Cliff(2, 0)$ [GTNut, Ch. VII.1, p. 410]. We need two gamma matrices that anticommute and square to one: $\gamma_1\gamma_2 = -\gamma_2\gamma_1$, $\gamma_1^2 = I$, and $\gamma_2^2 = I$, where I is the identity matrix. The first two Pauli matrices fit the bill:

$$\gamma_1 = \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \gamma_2 = \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}.$$

Then, we construct from those the generator of the spinor representation as $\frac{1}{2}\gamma_1\gamma_2$. The corresponding transformation is $U = e^{\gamma_1\gamma_2\theta/2}$:

$$\frac{1}{2}\gamma_1\gamma_2 = \frac{1}{2}i\sigma_3 = \frac{1}{2} \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix}, \quad U = \begin{pmatrix} e^{i\theta/2} & 0 \\ 0 & e^{-i\theta/2} \end{pmatrix}.$$

Now this looks like an honest spinor representation! But because $Spin(2) = SO(2)$, it is equivalent to the $S^+ \oplus S^-$ representation that we constructed earlier.