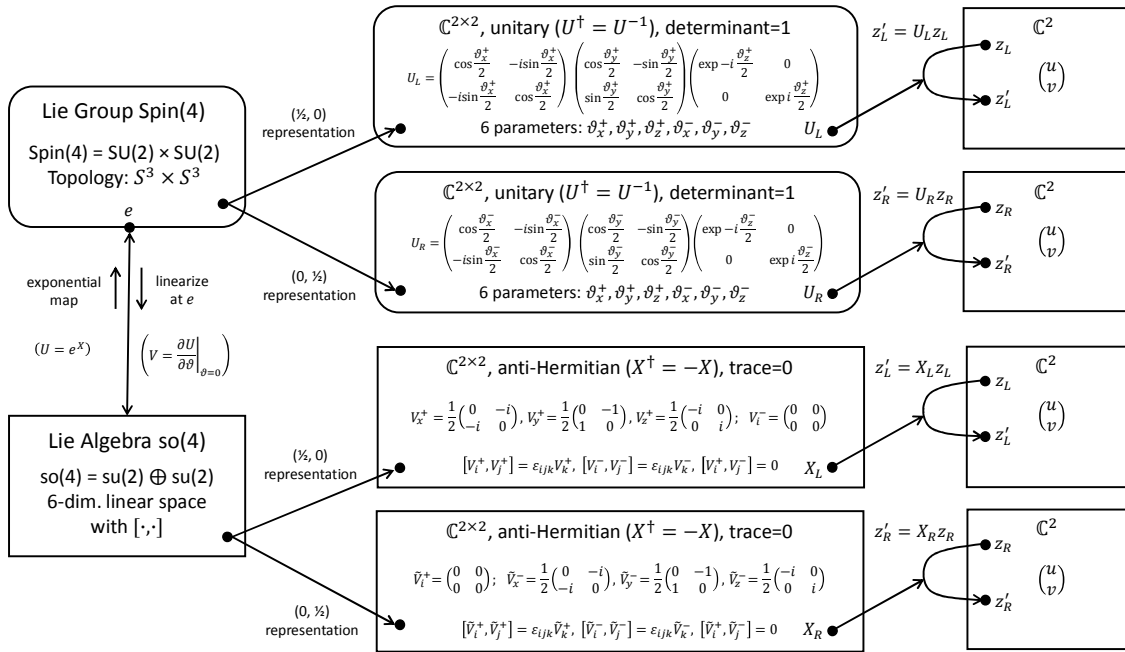


### 5.7 Spin(4): The (½, 0) and (0, ½) Representations



So far, we have used the 1- and 3-dimensional representations of so(3) to construct various representations of so(4). But we know that so(3) also has a 2-dimensional representation acting on complex 2-vectors known as *spinors*. Now, we are going to use this fact to construct 2-dimensional *spinor representations* of so(4)!

Let's start with the (½, 0) representation of so(4). Remembering that su(2) is isomorphic to so(3), we use the 2-dimensional (defining) representation (j = ½) of su(2), for the basis generators V<sub>i</sub><sup>+</sup>, that is,

$$V_x^+ = \frac{1}{2} \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix}, V_y^+ = \frac{1}{2} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, V_z^+ = \frac{1}{2} \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix},$$

and the 1-dimensional representation (j = 0) for the basis generators V<sub>i</sub><sup>-</sup>, that is,

$$V_x^- = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}, V_y^- = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}, V_z^- = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}.$$

These matrices satisfy all the commutation relations of so(4), including [V<sub>i</sub><sup>+</sup>, V<sub>j</sub><sup>-</sup>] = 0.

Exponentiating this Lie algebra yields the unitary transformation U<sub>L</sub> = exp(V<sub>x</sub><sup>+</sup>θ<sub>x</sub><sup>+</sup>) · exp(V<sub>y</sub><sup>+</sup>θ<sub>y</sub><sup>+</sup>) · exp(V<sub>z</sub><sup>+</sup>θ<sub>z</sub><sup>+</sup>), which we are already familiar with from SU(2). Note that this transformation depends only on the first three parameters, θ<sub>x</sub><sup>+</sup>, θ<sub>y</sub><sup>+</sup>, θ<sub>z</sub><sup>+</sup>; alternatively, if we go back to the plane-rotation parameters, it depends on θ<sub>x</sub>, θ<sub>y</sub>, θ<sub>z</sub> and φ<sub>x</sub>, φ<sub>y</sub>, φ<sub>z</sub> in the *same* way. We have found a 2-dimensional spinor representation of Spin(4) (see the upper branch of the diagram)!

The construction of the (0, ½) representation proceeds along the same lines. The main difference is that the unitary transformation U<sub>R</sub> now depends only on the second three parameters θ<sub>x</sub><sup>-</sup>, θ<sub>y</sub><sup>-</sup>, θ<sub>z</sub><sup>-</sup> or, if expressed in terms of plane-rotation parameters, on θ<sub>x</sub>, θ<sub>y</sub>, θ<sub>z</sub> and φ<sub>x</sub>, φ<sub>y</sub>, φ<sub>z</sub> in *opposite* ways. We

have found another 2-dimensional spinor representation of  $\text{Spin}(4)$  (see the lower branch of the diagram)!

Are the  $(\frac{1}{2}, 0)$  and  $(0, \frac{1}{2})$  representations equivalent? No, despite looking very similar! It is not possible to turn them into each other with a similarity transformation because the respective matrices depend on two different subsets of the six parameters. Remember that the  $(1, 0)$  and  $(0, 1)$  representations were also inequivalent. In fact, all  $(j, 0)$  and  $(0, j)$  representations of  $\text{Spin}(4)$  are inequivalent, irreducible representations.

Given the  $(\frac{1}{2}, 0)$  and  $(0, \frac{1}{2})$  representations, we can easily construct a (reducible) 4-dimensional representation of  $\text{Spin}(4)$  by taking the direct sum  $(\frac{1}{2}, 0) \oplus (0, \frac{1}{2})$ . Its basis generators are

$$V_x^+ = \frac{1}{2} \begin{pmatrix} 0 & -i & 0 & 0 \\ -i & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, V_y^+ = \frac{1}{2} \begin{pmatrix} 0 & -1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, V_z^+ = \frac{1}{2} \begin{pmatrix} -i & 0 & 0 & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix},$$

$$V_x^- = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & -i & 0 \end{pmatrix}, V_y^- = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{pmatrix}, V_z^- = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & -i & 0 \\ 0 & 0 & 0 & i \end{pmatrix}.$$

This representation is different from the *irreducible* 4-dimensional representation that defines  $\text{SO}(4)$ . We'll see how to construct the latter representation in a moment.